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So Hirata

Univ. of Illinois, Urbana, IL

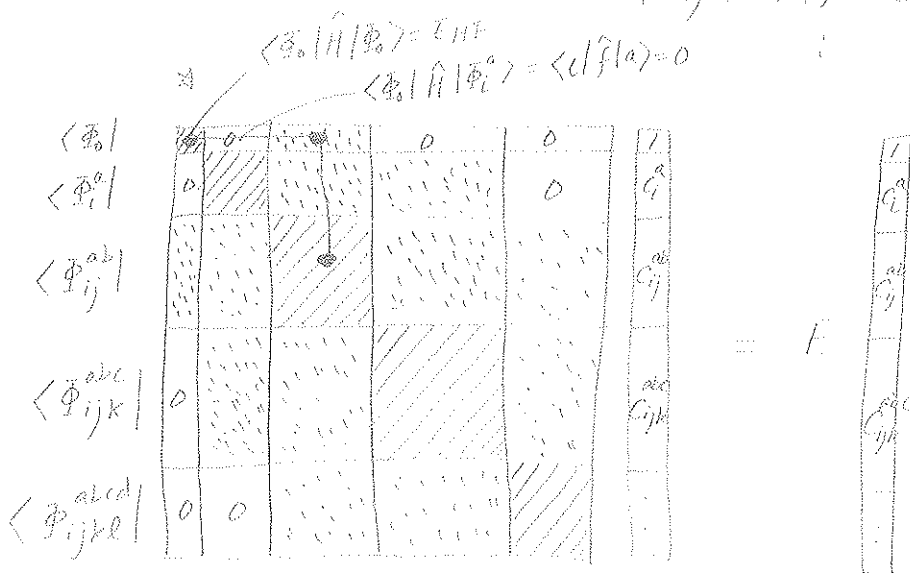
① Configuration interaction (CI) method

$$|\psi\rangle = |\Phi_0\rangle + \sum_{a,i} c_i^a |\Phi_i^a\rangle + \sum_{i,j} c_{ij}^{ab} |\Phi_{ij}^{ab}\rangle + \dots$$

$$\frac{\partial}{\partial c} \frac{\langle \psi | \hat{H} | \psi \rangle}{\langle \psi | \psi \rangle} = 0 \rightarrow$$

$$\begin{aligned} \langle \Phi_0 | \hat{H} | \psi \rangle &= E \langle \Phi_0 | \psi \rangle \\ \langle \Phi_i^a | \hat{H} | \psi \rangle &= E \langle \Phi_i^a | \psi \rangle \\ \langle \Phi_{ij}^{ab} | \hat{H} | \psi \rangle &= E \langle \Phi_{ij}^{ab} | \psi \rangle \end{aligned}$$

$\hat{H}|\psi\rangle = E|\psi\rangle$
 is true
 in the spaces of
 $\langle \Phi_0 |, \langle \Phi_i^a |, \langle \Phi_{ij}^{ab} |$



$$\langle \Phi_0 | \hat{H} | \Phi_i^a \rangle = \langle \Phi_i^a | \hat{H} | \Phi_0 \rangle^* = \langle i | \hat{f} | a \rangle = \langle a | \hat{f} | i \rangle^* = 0$$

(Brillouin theorem)

$$\langle \Phi_0 | \hat{H} | \Phi_{ijk}^{abc} \rangle = \langle \Phi_{ijk}^{abc} | \hat{H} | \Phi_0 \rangle^* = 0 \text{ (two-body nature of } \hat{H} \text{)}$$

CI should be cost-effective!

$$|\psi_{CID}\rangle = |\Phi_0\rangle + \sum_{\substack{i < j \\ a < b}} c_{ij}^{ab} |\Phi_{ij}^{ab}\rangle \text{ (truncated CI)}$$

→ "Intermediate" normalization ($c_0 = 1$)
 (in other words, $|\psi_{CID}\rangle$ is not normalized)

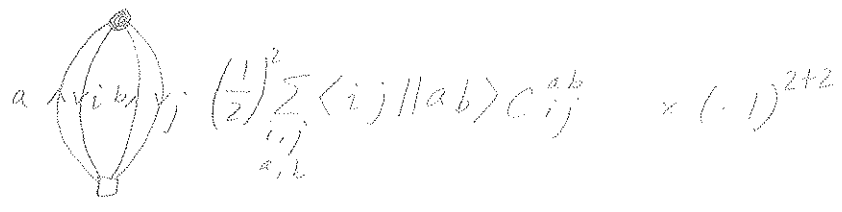
Energy eq.

$$E_{CID} = \langle \Phi_0 | \hat{H} | \Psi_{CID} \rangle = \underbrace{\langle \Phi_0 | \hat{H} | \Phi_0 \rangle}_{E_{HF}} + \underbrace{\sum_{\substack{ij \\ ab}} C_{ij}^{ab} \langle \Phi_0 | \hat{H} | \Phi_{ij}^{ab} \rangle}_{E_{corr.} (= E_{CID} - E_{HF})}$$

$$\left(\frac{1}{4}\right)^2 \sum_{\substack{ij \\ a,b}} C_{ij}^{ab} \langle p_8 || rs \rangle \langle \Phi_0 | \{ \hat{p}^{\dagger 1} \hat{q}^{\dagger 1} \hat{s}^{\dagger 1} \hat{r}^{\dagger 1} \} \{ \hat{a}^{\dagger 1} \hat{b}^{\dagger 1} \} | \Phi_0 \rangle$$

and 3 other contractions

OR



$$E_{corr.}^{CID} = \frac{1}{4} \sum_{\substack{ij \\ a,b}} \langle ij || ab \rangle C_{ij}^{ab}$$

C₂ amplitude eq.

$\frac{1}{4} \sum_{\substack{ij \\ cid}}$ missing

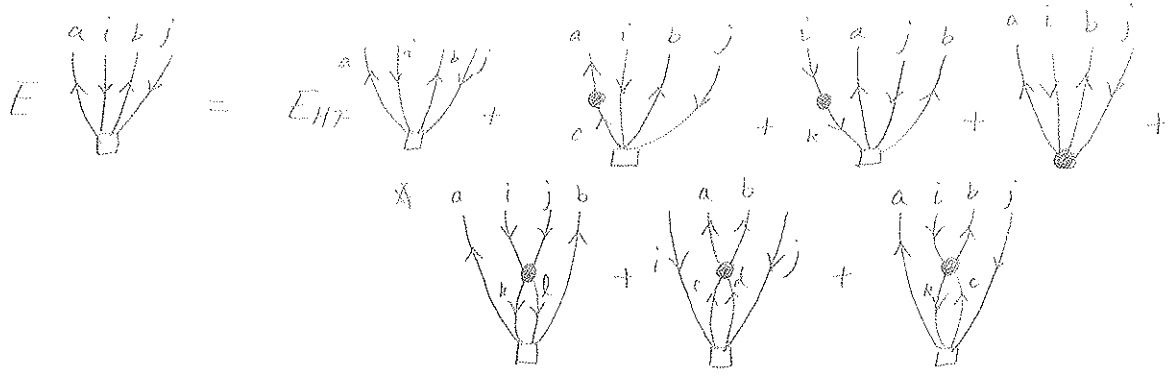
$$E C_{ij}^{ab} = \langle \Phi_{ij}^{ab} | \hat{H} | \Psi_{CID} \rangle = E_{HF} \langle \Phi_0 | \{ \hat{i}^{\dagger 1} \hat{j}^{\dagger 1} \} \{ \hat{a}^{\dagger 1} \hat{b}^{\dagger 1} \} \{ \hat{r}^{\dagger 1} \hat{s}^{\dagger 1} \} | \Phi_0 \rangle C_{cd}^{ab}$$

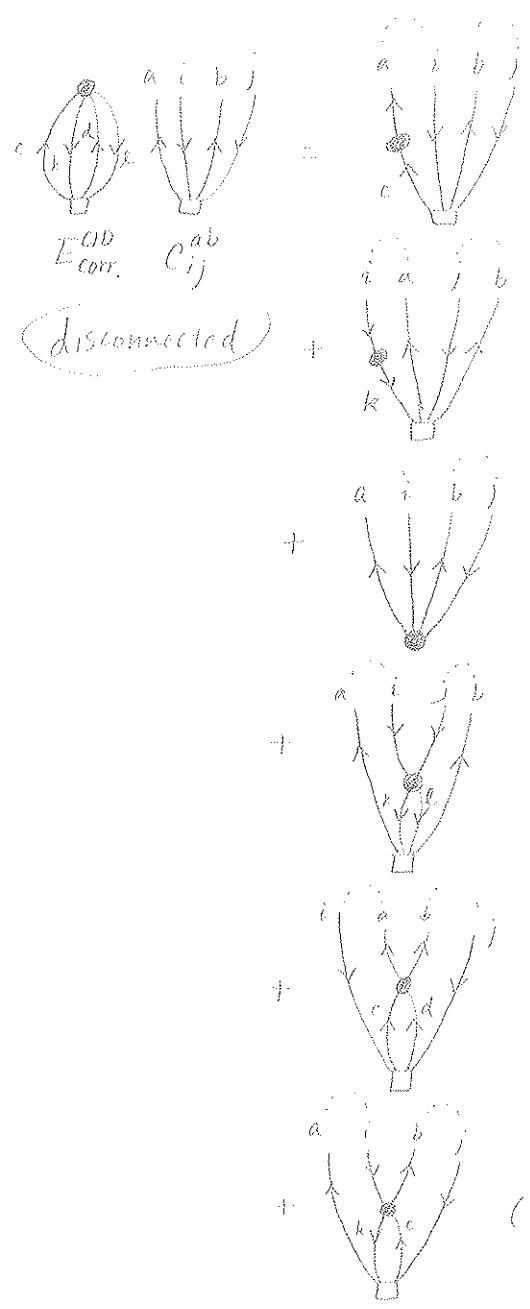
$$+ \sum_{p,q} \langle p || q \rangle \langle \Phi_0 | \{ \hat{i}^{\dagger 1} \hat{j}^{\dagger 1} \} \{ \hat{a}^{\dagger 1} \hat{b}^{\dagger 1} \} \{ \hat{p}^{\dagger 1} \hat{q}^{\dagger 1} \} \{ \hat{r}^{\dagger 1} \hat{s}^{\dagger 1} \} | \Phi_0 \rangle C_{cd}^{ab}$$

$$+ \frac{1}{4} \sum_{r,s} \langle p_8 || rs \rangle \langle \Phi_0 | \{ \hat{i}^{\dagger 1} \hat{j}^{\dagger 1} \} \{ \hat{a}^{\dagger 1} \hat{b}^{\dagger 1} \} \{ \hat{p}^{\dagger 1} \hat{q}^{\dagger 1} \} \{ \hat{s}^{\dagger 1} \hat{r}^{\dagger 1} \} | \Phi_0 \rangle$$

$$+ \frac{1}{4} \sum_{r,s} \langle p_8 || rs \rangle \langle \Phi_0 | \{ \hat{i}^{\dagger 1} \hat{j}^{\dagger 1} \} \{ \hat{a}^{\dagger 1} \hat{b}^{\dagger 1} \} \{ \hat{p}^{\dagger 1} \hat{q}^{\dagger 1} \} \{ \hat{s}^{\dagger 1} \hat{r}^{\dagger 1} \} \{ \hat{c}^{\dagger 1} \hat{d}^{\dagger 1} \} | \Phi_0 \rangle C_{cd}^{ab}$$

OR





$a-b$ non equivalent

$$(-1)^{2+2} P_{a|b} \sum_c \langle a|\hat{f}|c\rangle C_{ij}^{cb} = P_{a|b} \sum_c \langle a|\hat{f}|c\rangle C_{ij}^{cb} - \sum_c \langle b|\hat{f}|c\rangle C_{ij}^{ca} = (1 - \text{replace } a \leftrightarrow b) = (C_{a|b}) C_{ij}^{ab}$$

$$(-1)^{2+3} P_{i|j} \sum_k \langle k|\hat{f}|i\rangle C_{kj}^{ab} = -\sum_k \langle k|\hat{f}|i\rangle C_{kj}^{ab} + \sum_k \langle k|\hat{f}|j\rangle C_{ki}^{ab} = -(C_{i|j}) C_{ij}^{ab}$$

$$(-1)^{2+2} \langle ab||ij\rangle$$

$k-l$ equivalent

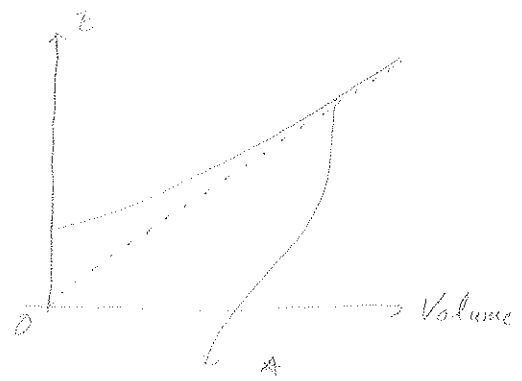
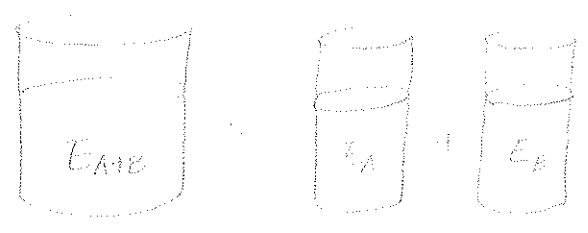
$$(-1)^{2+4} \begin{pmatrix} 1 \\ j \end{pmatrix} \sum_{k,l} \langle k||l||ij\rangle C_{kl}^{ab}$$

$c-d$

$$(-1)^{2+2} \begin{pmatrix} 1 \\ d \end{pmatrix} \sum_{c,d} \langle ab||cd\rangle C_{ij}^{cd}$$

$$(-1)^{2+3} P_{a|b} P_{i|j} \sum_{k,c} \langle kb||i|c\rangle C_{kj}^{ac}$$

Size consistency (size extensivity)



It's one of the foundations of statistical thermodynamics in chemistry!

Energy is "asymptotically" proportional to Volume

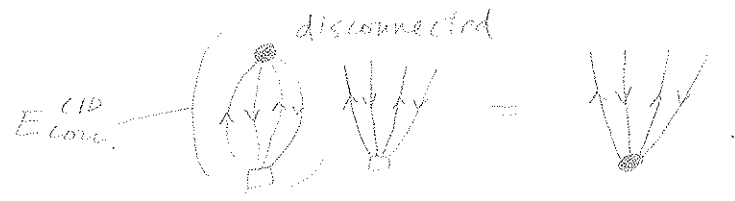
$$\lim_{V \rightarrow \infty} \frac{E}{V} = \text{finite}$$

$$\lim_{V \rightarrow \infty} E \propto V^1$$

Q: Is CID size consistent? *

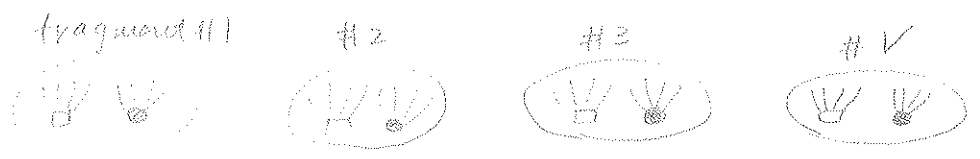
A: No! Thus, CID is a nonsensical method for total energy, (if not for excitations)

In the first approximation, CID eq. may be simplified to



This is what one gets in the first iteration of the iterative sol. of CI eq. with $C_{ij}^{ab} = 0$ initial guess.

In a local basis, we have local Ψ and Ψ for each fragment



$$E_{corr}^{CID} = V \int_{\text{local}} \Psi \Psi ; \text{ Assume } \int_{\text{local}} \Psi \Psi \propto V^n$$

$$\Psi = \frac{\int_{\text{local}} \Psi \Psi}{E_{corr}^{CID}} \propto \frac{V^n}{V^{n+1}} = V^{-n-1}$$

$$V^n \propto \int_{\text{local}} \Psi \Psi \propto V^{-n-1} \rightarrow n = -\frac{1}{2}, E_{corr}^{CID} = V \cdot V^{-\frac{1}{2}} = V^{\frac{1}{2}}$$

not extensive