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① Rayleigh-Schrödinger perturbation theory (RSPT)

$\hat{H} \psi_0 = \epsilon_0 \psi_0$ $\lambda = 1$

$\hat{H} = \hat{H}_0 + \lambda \hat{V}$

$\psi_0 = \psi_0^{(0)} + \lambda \psi_0^{(1)} + \lambda^2 \psi_0^{(2)} + \dots$

$\epsilon_0 = \epsilon_0^{(0)} + \lambda \epsilon_0^{(1)} + \lambda^2 \epsilon_0^{(2)} + \dots$

perturbation order

state (0 = GS, 1 = ES, ...)

$\hat{H}_0 \psi_k^{(0)} = \epsilon_k^{(0)} \psi_k^{(0)}$

$(\hat{H}_0 + \lambda \hat{V})(\psi_0^{(0)} + \lambda \psi_0^{(1)} + \lambda^2 \psi_0^{(2)} + \dots) = (\epsilon_0^{(0)} + \lambda \epsilon_0^{(1)} + \lambda^2 \epsilon_0^{(2)} + \dots)(\psi_0^{(0)} + \lambda \psi_0^{(1)} + \lambda^2 \psi_0^{(2)} + \dots)$

$\lambda^0: \hat{H}_0 \psi_0^{(0)} = \epsilon_0^{(0)} \psi_0^{(0)}$ (A)

$\lambda^1: \hat{H}_0 \psi_0^{(1)} + \hat{V} \psi_0^{(0)} = \epsilon_0^{(1)} \psi_0^{(1)} + \epsilon_0^{(1)} \psi_0^{(0)}$ (B) *

$\lambda^2: \hat{H}_0 \psi_0^{(2)} + \hat{V} \psi_0^{(1)} = \epsilon_0^{(2)} \psi_0^{(2)} + \epsilon_0^{(1)} \psi_0^{(1)} + \epsilon_0^{(2)} \psi_0^{(0)}$ (C)

$\lambda^3: \hat{H}_0 \psi_0^{(3)} + \hat{V} \psi_0^{(2)} = \epsilon_0^{(3)} \psi_0^{(3)} + \epsilon_0^{(1)} \psi_0^{(2)} + \epsilon_0^{(2)} \psi_0^{(1)} + \epsilon_0^{(3)} \psi_0^{(0)}$ (D)

(E) $\psi_0^{(1)} = \sum_{k \neq 0} c_k^{(1)} \psi_k^{(0)}$

$\sum_{k \neq 0} c_k^{(1)} \epsilon_k^{(0)} \psi_k^{(0)} + \hat{V} \psi_0^{(0)} = \epsilon_0^{(1)} \sum_{k \neq 0} c_k^{(1)} \psi_k^{(0)} + \epsilon_0^{(1)} \psi_0^{(0)}$ $\int \psi_0^{(0)*} d\tau$

$\int \psi_0^{(0)*} d\tau \left(\epsilon_0^{(1)} = \int \psi_0^{(0)*} \hat{V} \psi_0^{(0)} d\tau \right)$

$c_k^{(1)} = \frac{\int \psi_k^{(0)*} \hat{V} \psi_0^{(0)} d\tau}{\epsilon_0^{(0)} - \epsilon_k^{(0)}}$

(F) $\psi_0^{(2)} = \sum_{k \neq 0} c_k^{(2)} \psi_k^{(0)}$ *

$\sum_{k \neq 0} c_k^{(2)} \epsilon_k^{(0)} \psi_k^{(0)} + \hat{V} \sum_{k \neq 0} c_k^{(1)} \psi_k^{(0)} = \epsilon_0^{(2)} \sum_{k \neq 0} c_k^{(2)} \psi_k^{(0)} + \epsilon_0^{(1)} \sum_{k \neq 0} c_k^{(1)} \psi_k^{(0)} + \epsilon_0^{(2)} \psi_0^{(0)}$

$\int \psi_k^{(0)*} d\tau \left(\epsilon_0^{(2)} = \sum_{k \neq 0} c_k^{(1)} \int \psi_0^{(0)*} \hat{V} \psi_k^{(0)} d\tau = \sum_{k \neq 0} \frac{\langle \psi_0^{(0)} | \hat{V} | \psi_k^{(0)} \rangle \langle \psi_k^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{\epsilon_0^{(0)} - \epsilon_k^{(0)}} \int \psi_0^{(0)*} d\tau \right)$

$c_k^{(2)} \epsilon_k^{(0)} + \sum_{k' \neq 0} c_{k'}^{(1)} \int \psi_k^{(0)*} \hat{V} \psi_{k'}^{(0)} d\tau = c_k^{(2)} \epsilon_0^{(0)} + c_k^{(1)} \epsilon_0^{(1)}$

$c_k^{(2)} = \sum_{k' \neq 0} \frac{\langle \psi_k^{(0)} | \hat{V} | \psi_{k'}^{(0)} \rangle \langle \psi_{k'}^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{(\epsilon_0^{(0)} - \epsilon_k^{(0)})(\epsilon_0^{(0)} - \epsilon_{k'}^{(0)})} - \epsilon_0^{(1)} \frac{\langle \psi_k^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{(\epsilon_0^{(0)} - \epsilon_k^{(0)})^2}$ *

(V)

$\epsilon_0^{(2)} = \langle \psi_0^{(0)} | \hat{V} | \psi_0^{(0)} \rangle = \sum_{k, k' \neq 0} \frac{\langle \psi_0^{(0)} | \hat{V} | \psi_k^{(0)} \rangle \langle \psi_k^{(0)} | \hat{V} | \psi_{k'}^{(0)} \rangle \langle \psi_{k'}^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{(\epsilon_0^{(0)} - \epsilon_k^{(0)})(\epsilon_0^{(0)} - \epsilon_{k'}^{(0)})} - \epsilon_0^{(1)} \sum_{k \neq 0} \frac{\langle \psi_0^{(0)} | \hat{V} | \psi_k^{(0)} \rangle \langle \psi_k^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{(\epsilon_0^{(0)} - \epsilon_k^{(0)})^2}$

② Møller-Plesset (MP) partitioning

$$\hat{H} = I_{HF} + \left(\sum_{p,q} \langle p|\hat{f}|q\rangle \{\hat{p}^\dagger \hat{q}\} + \frac{1}{4} \sum_{p,q,r,s} \langle pq||rs\rangle \{\hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r}\} \right)$$

$\left(\sum_i \langle i|\hat{f}|i\rangle - \frac{1}{2} \sum_{ij} \langle ij||ij\rangle \right) \hat{H}_0$ (rest is \hat{V})

$$E_{MP}^{(0)} = \langle \psi_0^{(0)} | \hat{H}_0 | \psi_0^{(0)} \rangle = \sum_i \langle i|\hat{f}|i\rangle + \underbrace{\sum_{p,q} \langle p|\hat{f}|q\rangle \langle \psi_0^{(0)} | \{\hat{p}^\dagger \hat{q}\} | \psi_0^{(0)} \rangle}_0$$

$$= \sum_i \langle i|\hat{f}|i\rangle$$

$$E_{MP}^{(1)} = \langle \psi_0^{(0)} | \hat{V} | \psi_0^{(0)} \rangle = -\frac{1}{2} \sum_{ij} \langle ij||ij\rangle + \underbrace{\langle \psi_0^{(0)} | \{\hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r}\} | \psi_0^{(0)} \rangle}_0$$

$E_{MP}^{(0)} + E_{MP}^{(1)} = \langle \psi_0^{(0)} | \hat{H} | \psi_0^{(0)} \rangle = E_{HF}$

$$E_{MP}^{(2)} = \sum_{k \neq 0} \frac{\langle \psi_0^{(0)} | \hat{V} | \psi_k^{(0)} \rangle \langle \psi_k^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{E_0^{(0)} - E_k^{(0)}}$$

$$E_0^{(0)} = \sum_i \langle i|\hat{f}|i\rangle$$

$$E_k^{(0)} = \sum_i \langle i|\hat{f}|i\rangle + \sum_{p,q} \langle p|\hat{f}|q\rangle$$

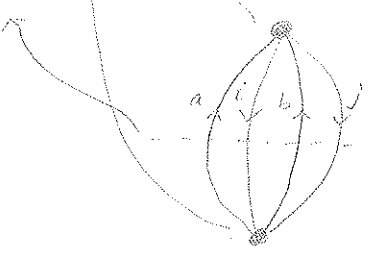
$$\frac{1}{4} \sum_{p,q,r,s} \langle pq||rs\rangle \langle \Phi_0 | \{\hat{p}^\dagger \hat{q}^\dagger \hat{s} \hat{r}\} \{ \hat{a}^\dagger \hat{b}^\dagger \hat{c} \hat{d} \} | \Phi_0 \rangle$$

$$= \langle ij||ab\rangle \quad (\text{const. part contributes nothing})$$

$$\langle \Phi_0 | \hat{V} | \psi_k^{(0)} \rangle = \sum_i \langle i|\hat{f}|i\rangle + \underbrace{\langle a|\hat{f}|a\rangle}_{E_a} + \underbrace{\langle b|\hat{f}|b\rangle}_{E_b} - \underbrace{\langle i|\hat{f}|i\rangle}_{E_i} - \underbrace{\langle j|\hat{f}|j\rangle}_{E_j}$$

$$\frac{1}{4} \sum_{ij,a,b} \frac{\langle ij||ab\rangle \langle ab||ij\rangle}{E_i + E_j - E_a - E_b}$$

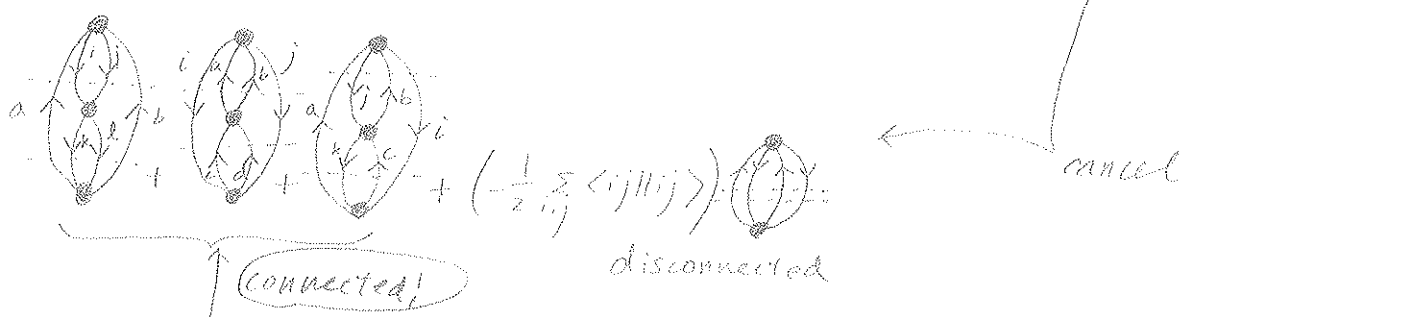
two equivalent pairs



(1) 2+2

$$E_{MP}^{(2)} = \sum_{k, k+0} \frac{\langle \psi_0^{(0)} | \hat{V} | \psi_k^{(0)} \rangle \langle \psi_k^{(0)} | \hat{V} | \psi_{k'}^{(0)} \rangle \langle \psi_{k'}^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{(E_0 - E_k)(E_0 - E_{k'})} - \frac{\langle \psi_0^{(0)} | \hat{V} | \psi_k^{(0)} \rangle \langle \psi_k^{(0)} | \hat{V} | \psi_0^{(0)} \rangle}{(E_0 - E_k)^2}$$

Diagrammatic expansion of the second-order energy correction. The first term is a sum over two intermediate states \$k\$ and \$k'\$, represented by a diagram with three vertices and two internal lines. The second term is a sum over a single intermediate state \$k\$, represented by a diagram with two vertices and one internal line. The energy denominators are \$(E_0 - E_k)(E_0 - E_{k'})\$ and \$(E_0 - E_k)^2\$ respectively.



$$E_{MP}^{(2)} = (-1)^{2+4} \left(\frac{1}{2}\right)^3 \sum_{\substack{i,j \\ a,b \\ k,l}} \frac{\langle ij || ab \rangle \langle ab || kl \rangle \langle kl || ij \rangle}{(e_i + e_j - e_a - e_b)(e_k + e_l - e_c - e_d)}$$

$$+ (-1)^{2+2} \left(\frac{1}{2}\right)^3 \sum_{\substack{i,j \\ a,b \\ c,d}} \frac{\langle ij || ab \rangle \langle ab || cd \rangle \langle cd || ij \rangle}{(e_i + e_j - e_a - e_b)(e_c + e_d - e_c - e_d)}$$

$$+ (-1)^{2+3} \sum_{\substack{i,j,k \\ a,b,c}} \frac{\langle ij || ab \rangle \langle kb || jc \rangle \langle ac || ki \rangle}{(e_i + e_j - e_a - e_b)(e_i + e_k - e_a - e_c)}$$

③ CCD redux

$$E_{corr.}^{(CD)} = \text{diagram} = \frac{1}{4} \sum_{\substack{ij \\ a,b}} \langle ij || ab \rangle T_{ij}^{ab} \quad \text{--- (A)}$$

$$0 = \frac{(e_a + e_b - e_i - e_j) T_{ij}^{ab}}{\text{diagram}} + \langle ab || ij \rangle \text{diagram} + \dots$$

connected!

all connected

Consider solving (B) iteratively starting w/ $T_{ij}^{ab} = 0$ except the first two terms

$$T_{ij}^{ab[1]} = \frac{\langle ab || ij \rangle}{(e_i + e_j - e_a - e_b)} = \text{diagram}$$

$$E_{corr.}^{(CD)[1]} = \frac{1}{4} \sum_{\substack{ij \\ a,b}} \langle ij || ab \rangle T_{ij}^{ab[1]} = \frac{1}{4} \sum_{\substack{ij \\ a,b}} \frac{\langle ij || ab \rangle \langle ab || ij \rangle}{(e_i + e_j - e_a - e_b)} = \text{diagram} = E_{MP}^{(2)}$$

$$T_{ij}^{ab[2]} = \text{diagram} + \left[\text{diagram} + \text{diagram} + \text{diagram} \right] / (e_i + e_j - e_a - e_b)$$

$$= \text{diagram} + \text{diagram} + \text{diagram} + \text{diagram}$$

$$E_{corr.}^{(CD)[2]} = \text{diagram} + \text{diagram} + \text{diagram} + \text{diagram} = E_{MP}^{(2)} + E_{MP}^{(2)}$$

connected!

④ Goldstone's linked-diagram theorem

$$E_{MP}^{(n)} = \left. \begin{matrix} \text{diagram} & 1 \\ \text{diagram} & 2 \\ \vdots & \vdots \\ \text{diagram} & n-1 \\ \text{diagram} & n \end{matrix} \right\} \text{connected}$$

March et al. time-dependent derivation
Shavitt & Bartlett time-independent derivation